

# 10/7 Squeezed States Lecture 14

$$i\hbar \frac{\partial}{\partial t} \psi(x,t) = \frac{-\hbar^2}{2m(t)} \frac{\partial^2}{\partial x^2} \psi(x,t) + \frac{1}{2} m(t) \omega_0^2 x^2 \psi(x,t) - F(t) x \psi(x,t) + v(t) i\hbar \frac{\partial}{\partial x} \psi(x,t)$$

$$\psi(x,t) = \left( \frac{2\alpha(t)}{\pi} \right)^{1/4} e^{-i\theta(t)} e^{i p(t)(x-\bar{x}(t)) / \hbar} e^{-\alpha(t) (x-\bar{x}(t))^2} v(t) \hat{p} \psi$$

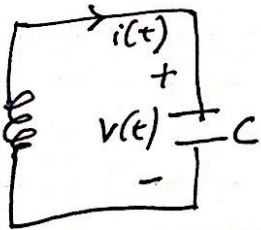
$$\frac{d}{dt} \bar{x}(t) = \frac{p(t)}{m(t)} - v(t)$$

$$\frac{d}{dt} p(t) = -m(t) \omega_0^2(t) \bar{x}(t) + F(t) - F(t) \bar{x}(t)$$

$$\hbar \frac{d}{dt} \theta(t) = \frac{\hbar^2 \alpha(t)}{2m(t)} + \frac{m(t) \omega_0^2(t)}{8\alpha(t)} - \frac{p^2(t)}{2m(t)} + \frac{1}{2} m(t) \omega_0^2(t) \bar{x}(t)$$

$$i\hbar \frac{d}{dt} \alpha(t) = \frac{2\hbar^2 \alpha^2(t)}{m(t)} - \frac{1}{2} m(t) \omega_0^2(t)$$

## LC-Circuit



$$\frac{d}{dt} v(t) = \frac{1}{C} i(t)$$

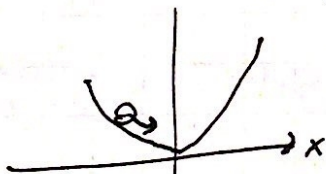
$$\frac{d}{dt} x(t) = \frac{p(t)}{m}$$

$$\frac{d}{dt} i(t) = -\frac{1}{L} v(t)$$

$$\frac{d}{dt} p(t) = -m\omega_0^2 x(t)$$

$$E = \frac{1}{2} L i^2(t) + \frac{1}{2} (v^2(t))$$

$$E = \frac{p^2(t)}{2m} + \frac{1}{2} m\omega_0^2 x^2(t)$$



$$V(x) = \frac{1}{2} m \omega_0^2 x^2$$

$$E = \frac{p^2(t)}{2m} + \frac{1}{2} m\omega_0^2 x(t)$$

$$v(t) \leftrightarrow x(t)$$

$$i(t) \leftrightarrow p(t)$$

$$\hat{E} \psi(x,t) = \frac{\hat{p}^2}{2m} \psi(x,t) + \frac{1}{2} m\omega_0^2 x^2 \psi(x,t)$$

$$i\hbar \frac{\partial}{\partial t} \psi(x,t) = \frac{(-i\hbar \frac{\partial}{\partial x})^2}{2m} \psi(x,t) + \frac{1}{2} m\omega_0^2 x^2 \psi(x,t)$$

$$E = \frac{1}{2} L i^2(t) + \frac{1}{2} (v^2(t))$$

$$\hat{E} \psi(v,t) = \frac{1}{2} L i^2 \psi(v,t) + \frac{1}{2} (v^2) \psi(v,t)$$

$$i\hbar \frac{\partial}{\partial t} \psi(v,t) = \frac{1}{2} L (-iA \frac{\partial}{\partial v})^2 \psi(v,t) + \frac{1}{2} (v^2) \psi(v,t)$$

$$A = \omega_0^2 = \frac{1}{LC}$$

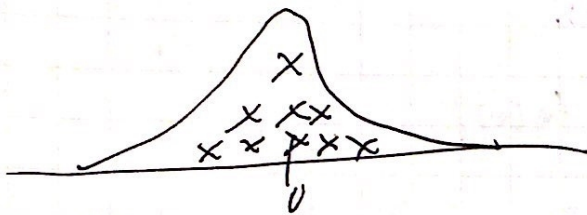
$$\hat{H} = -\frac{\hbar^2 \omega_0^2}{2L} \frac{\partial^2}{\partial v^2} + \frac{1}{2} C v^2$$

quantum mechanical  
sol'n for LC circuit.

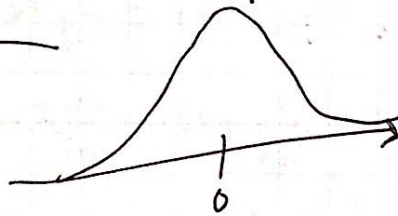
$$E \psi(v) = -\frac{\hbar^2 \omega_0^2}{2L} \frac{d^2}{dv^2} \psi(v) + \frac{1}{2} C v^2 \psi(v)$$

$$\psi_n(v) = \left[ \frac{C}{\pi \hbar \omega_0} \right]^{1/4} \frac{1}{\sqrt{2^n n!}} e^{-\frac{1}{2} \frac{C v^2}{\hbar \omega_0}} H_n \left( \sqrt{\frac{C}{\hbar \omega_0}} v \right)$$

$$E_n = \hbar \omega_0 \left( n + \frac{1}{2} \right)$$



$$\rho(v) = |\psi_0(v)|^2$$



Ground State

$$\int_{-\infty}^{\infty} \psi^*(v) \left( E \psi(v) = \frac{1}{2} L \hat{p}^2 \psi(v) + \frac{1}{2} C v^2 \psi(v) \right) dv$$

$$E = \frac{1}{2} L \langle \hat{p}^2 \rangle + \frac{1}{2} C \langle v^2 \rangle = \frac{1}{2} \hbar \omega_0$$

$$\langle v^2 \rangle = \langle v \rangle^2 + (\Delta v)^2$$

$$\langle \hat{p}^2 \rangle = \langle \hat{p} \rangle^2 + (\Delta p)^2$$

$$= \frac{1}{2} L (\Delta p)^2 + \frac{1}{2} C (\Delta v)^2$$

$$\frac{1}{2} C (\Delta v)^2 = \frac{1}{4} \hbar \omega_0 \Rightarrow (\Delta v)^2 = \frac{1}{2} \frac{\hbar \omega_0}{C}$$

$$\frac{1}{2} L (\Delta p)^2$$

$$= \frac{1}{4} \hbar \omega_0$$

$$= (\Delta p)^2 = \frac{1}{2} \frac{\hbar \omega_0}{L}$$

$$\Delta v \Delta p = \sqrt{\frac{1}{2} \frac{\hbar \omega_0}{C}} \cdot \sqrt{\frac{1}{2} \frac{\hbar \omega_0}{L}} = \frac{1}{2} \frac{\hbar \omega_0}{\sqrt{LC}}$$

$$\Delta v \Delta p \geq \frac{1}{2} \frac{\hbar \omega_0}{\sqrt{LC}}$$

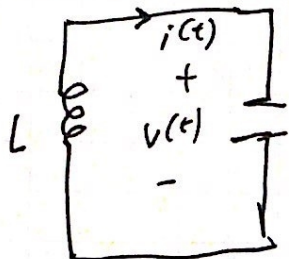
Gaussian  
ground state



$$\Rightarrow V = \sqrt{\frac{\hbar \omega_0}{2L}} (\hat{a} + \hat{a}^\dagger)$$

$$\hat{I} = \sqrt{\frac{\hbar \omega_0}{2L}} \left( \frac{\hat{a} - \hat{a}^\dagger}{i} \right)$$

$$\langle \phi_n(V) | V | \phi_n(V) \rangle = \sqrt{\frac{\hbar \omega_0}{2L}} (\sqrt{n} \int_{n', n-1} + \sqrt{n+1} \int_{n', n+1})$$



$$V(t) = -L \frac{d}{dt} i(t)$$

From pset

$$i(t) = \frac{d}{dt} q(t)$$

\* Don't know grand state energy

$$V(t) = \frac{d}{dq} E_q(q(t))$$

linear inductor,  
non-linear capacitor

$$E = \frac{1}{2} L i^2(t) + E_q(q(t))$$

$$\frac{d}{dt} q(t) = i(t) \quad \frac{d}{dt} x(t) = \frac{p(t)}{m}$$

$$\frac{d}{dt} i(t) = -\frac{1}{L} \frac{d}{dq} E_q(q(t))$$

$$\frac{d}{dt} p(t) = -\frac{d}{dx} V(x(t))$$

$$E = \frac{1}{2} L i^2(t) + E_q(q(t)) \quad E = \frac{p^2(t)}{2m} + V(x(t))$$

$$q(t) \leftrightarrow x(t)$$

$$i(t) \leftrightarrow p(t)$$

doesn't matter to  
choose for  $\hat{I}$

$$\hat{I} = -i\hbar \frac{\partial}{\partial q}$$

$$E = \frac{1}{2} L i^2(t) + E_q(q(t))$$

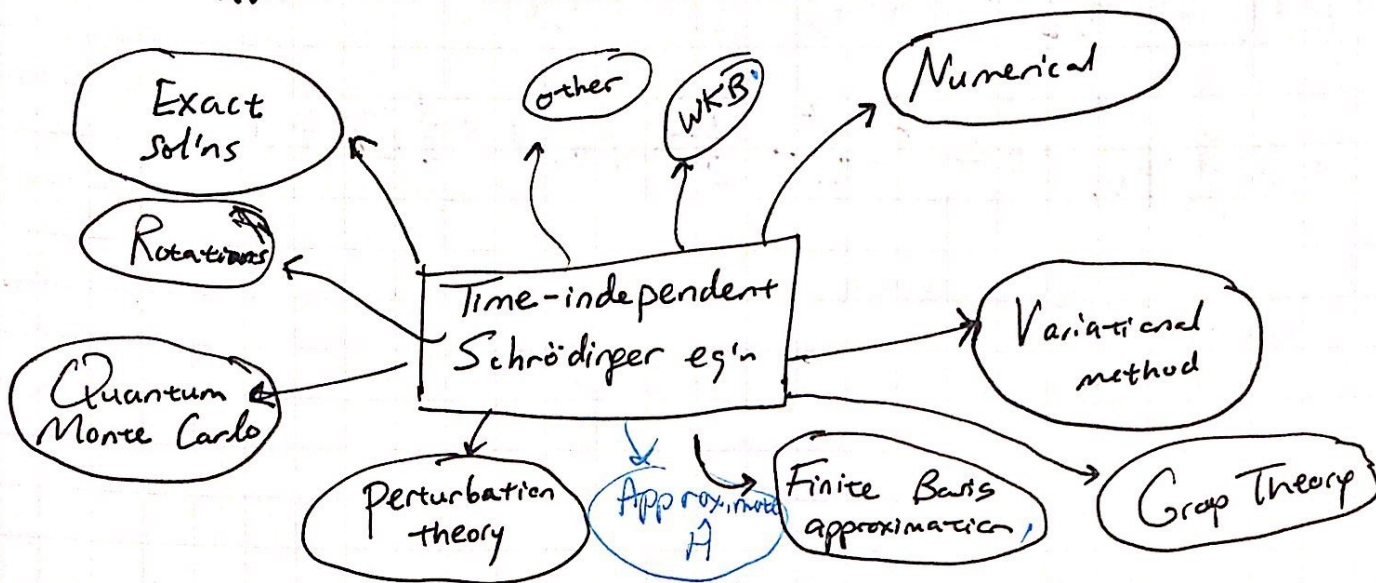
$$\hat{E} \Psi(q, t) = \frac{1}{2} L \hat{I}^2 \Psi(q, t) + E_q(q) \Psi(q, t)$$

$$i\hbar \frac{\partial}{\partial t} \Psi(q, t) = \frac{1}{2} L \left( i\hbar \frac{\partial}{\partial q} \right)^2 \Psi(q, t) + E_q(q) \Psi(q, t)$$

Use  
Ehrenfest's  
theorem:  $\frac{d}{dt} \langle q \rangle = \langle \hat{I} \rangle$

$$\frac{d}{dt} \langle \hat{I} \rangle = -\frac{1}{L} \left\langle \frac{d}{dq} E_q \right\rangle$$

4 models for which there are exact sol'n's



WKB : Jettieries, Wenzel, Kramers, Brillouin

Office hours 4 o'clock.